Estimation of dynamic buckling loads in approximate shell models via frequency-domain methods

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We study the stability and dissipativity of viscoelastic plates with constant thickness under dynamic loads given by the equations ([1, 7])

$$\frac{D}{h}(1-R^*)\nabla^4 w = L(w,F) - \rho \frac{\partial^2 w}{\partial t^2} ,
\frac{1}{E}\nabla^4 F = -\frac{1}{2}(1-R^*)L(w,w) .$$
(1)

In system (1) w is the deflection, F is the stress function, D is the plate's boundary stiffness, h is the thickness, ρ is the material's density and E is the elastic modulus. R^* denotes an integral operator with the relaxation kernel $R(t) = kt^{\alpha-1}e^{-\varepsilon t}$, where k > 0, $\varepsilon > 0$ and $\alpha \in [0, 1]$ are parameters. The expression L(w, w) is defined by

$$L(w,w) = 2 \frac{\partial^2 w}{\partial x_1^2} \frac{\partial^2 w}{\partial x_2^2} - 2 \left(\frac{\partial^2 w}{\partial x_1 \partial x_2}\right)^2.$$
(2)

The boundary conditions may be expressed as:

$$w = \frac{\partial^2 w}{\partial x_1^2} + \mu \frac{\partial^2 w}{\partial x_2^2} = 0 \quad \text{at} \quad x_1 = 0 \quad \text{and} \quad x_1 = a ;$$
(3)

$$w = \frac{\partial^2 w}{\partial x_2^2} + \mu \frac{\partial^2 w}{\partial x_1^2} = 0 \quad \text{at} \quad x_2 = 0 \quad \text{and} \quad x_2 = b;$$
(4)

$$\frac{\partial^2 F}{\partial x_2^2} = -\frac{1}{h} \left[k_1 + k_2 \cos \theta t \right] \quad \text{and} \quad \frac{\partial^2 F}{\partial x_1 \partial x_2} = 0 \quad \text{at} \quad x_1 = 0 \quad \text{and} \quad x_1 = a ; \qquad (5)$$

$$\frac{\partial^2 F}{\partial x_1^2} = 0$$
 and $\frac{\partial^2 F}{\partial x_1 \partial x_2} = 0$ at $x_2 = 0$ and $x_2 = b$. (6)

In (3) – (6) μ is Poisson's ratio ; a > 0 and b > 0 are the length and the width of the plate ; k_1, k_2 and θ are certain real parameters.

If we assume that the plate is hinged at its boundaries, the solution w of the boundaryvalue problem (1) - (6) can be expressed as

$$w(x_1, x_2, t) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} w_{mn}(t) \sin \frac{m\pi x_1}{a} \sin \frac{n\pi x_2}{b}.$$
 (7)

Using the Bubnov-Galerkin method one obtains ([1, 7]) for a = b = l a system of nonlinear integro-differential equations which can be written for $z := w_{mn}$ as

$$\ddot{z} + \Omega^2 (1 - q_0 \cos \theta t) z - \omega^2 R^* z + q_1 z (1 - R^*) z^2 = 0, \qquad (8)$$

where $(R^*f)(t) = k_1 \int_0^t k \tau^{\alpha-1} e^{-\varepsilon \tau} f(\tau) d\tau$ and $\Omega, \omega, q_0, q_1, k_1$ and θ are also parameters. For certain cases of these parameters we get the linear Mathieu equation

$$\ddot{z} + \Omega^2 (1 - q_0 \cos \theta t) z = 0 \tag{9}$$

and, with a further parameter q_4 , the nonlinear Mathieu-Duffing equation

$$\ddot{z} + \Omega^2 (1 - q_0 \cos \theta t) z + q_4 z^3 = 0.$$
(10)

Exploring frequency-domain methods developed in [10, 5] we derive sufficient conditions for the boundedness of solutions on \mathbb{R}_+ and the dissipativity of equations of the type (8) for nonsingular kernels ($\alpha = 1$). Under these conditions dynamic buckling is impossible. In particular we investigate the class of systems with periodic in time linear and nonlinear parts with the help of the frequency theorems for nonautonomous periodic in time equations ([9, 3]). We also consider the case of a singular kernel ($\alpha \in (0, 1)$). This leads to ordinary differential equations with fractional derivatives ([2]). The functional-analytic investigation of such singular differential equations can be done in certain interpolation spaces ([8]). For the characterization of boundedness properties of the solutions and their regularity we use the solvability of Riccati equations for plate problems ([6]).

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